# ELASTIC ANALYSIS OF THIN FIBER-REINFORCED PLATES 

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#### Abstract

This paper is concerned with the bending problem of a fibre-composite plate in the elastic range. Within the classical Kirchhoff plate theory, two approaches are utilized. In the model due to Świtka, the fibre-composite plate is treated as a homogeneous medium with a number of embedded families of fibres, while in the second model the plate is considered as a system of orthotropic layers. The bending problem is formulated in classical and variational forms and solved by finite element method. We have developed our own computer code and numerically compared predictions of the two models and other models available in the program ABAQUS on test examples. Numerical results illustrating influence of the placement of fibres on the quantities of interest are also included.


Key words: fibre-composite layered plate, bending of anisotropic plate, finite element method, elastic solution

## 1. INTRODUCTION

A composite material is formed when two or more materials are combined. The properties of the composite are different, usually better then its individual components. This paper deals with the problem of bending of laminated composite plates, including fibre-composite plates. The analysed composite plate consists of the matrix, which is made of a modern plastic material, and a set of regularly disposed fibres made of glass or titan, for example. The bending problem of composite plates has been studied in many works, see e.g. [2, 3, 4, 6]. A finite element method formulation of the model suggested in [4] is given in [5].

The properties of a composite dependent on the properties of the constituent materials, and their distribution and interaction. In this paper we have
treated the composite plate as a fibre-composite plate (model of Świtka) and as a system of layers (a layerwise model). We have based our analysis on the basic assumptions of classical Kirchhoff plate theory.

Our aim in this is twofold; firstly, to develop our own computer program for anisotropic plates and to test it on typical examples, secondly we want to analyse the influence of the fibre orientation on the plate deflection and the plate bending and twisting moments. In the sequel, we shall briefly present two models for an anisotropic plate, constitutive relations for an orthotropic layer and next formulate the bending problem in classical and variational forms and its FEM approximation. We have developed our own computer code and numerically compared predictions of the model by Świtka and the layerwise model and other models available in the program ABAQUS on test examples. Numerical results illustrating influence of the placement of fibres on the quantities of interest are also included. The obtained results

## 2. THIN ANISOTROPIC PLATE

The differential equation for a thin anisotropic plate in bending takes the form

$$
\begin{gather*}
D_{11} \frac{\partial^{4} w}{\partial x^{4}}+\left(2 D_{12}+4 D_{66}\right) \frac{\partial^{4} w}{\partial x^{2} \partial y^{2}}+D_{22} \frac{\partial^{4} w}{\partial y^{4}}+ \\
4 D_{16} \frac{\partial^{4} w}{\partial x^{3} \partial y}+4 D_{26} \frac{\partial^{4} w}{\partial x \partial y^{3}}=p \tag{1}
\end{gather*}
$$

wherein $w=w(x, y)$ is the displacement of the middle plane of the plate $\Omega \subset R^{2}, \quad p=p(x, y)$ is the load on the plate, and coefficients $D_{i j}$ represent the plate bending and twisting stiffness,

$$
\mathbf{D}=\left[\begin{array}{lll}
D_{11} & D_{12} & D_{16}  \tag{2}\\
D_{12} & D_{22} & D_{26} \\
D_{16} & D_{26} & D_{66}
\end{array}\right]
$$

### 2.1. The model of Świtka

The cross-section of the analysed fibre-composite plate is shown in fig. 1 in which, for clarity, only one $r$ th family of fibres is shown.


Fig.1. Fibre composite plate with one family of fibres
The differential equation for a fibre-composite plate in bending takes the form

$$
\begin{equation*}
D \nabla^{4} w+\left[\sum_{r} \frac{E_{r} A_{r}}{b_{r}} z_{r}^{2} s_{i}^{r} s_{j}^{r} s_{k}^{r} s_{l}^{r}-\sum_{r} \frac{E_{r} A_{r}}{b_{r}} z_{r} s_{i}^{r} s_{j}^{r} s_{p}^{r} s_{q}^{r} \sum_{s} c_{p q}^{s} z_{s} s_{k}^{s} s_{l}^{s}\right] w_{, i j k l}=p \tag{3}
\end{equation*}
$$

Having compared coefficients of eqns. (1) and (3) we arrive at the following relationships (4):

$$
\begin{align*}
D_{11}= & D+\sum_{r} \frac{E_{r} A_{r}}{b_{r}} z_{r}^{2}\left(s_{1}^{r}\right)^{4}-\sum_{r} \frac{E_{r} A_{r}}{b_{r}} z_{r}\left(s_{1}^{r}\right)^{2} s_{p}^{r} s_{q}^{r} \sum_{s} c_{p q}^{s} z_{s}\left(s_{l}^{s}\right)^{2} \\
D_{22}= & D+\sum_{r} \frac{E_{r} A_{r}}{b_{r}} z_{r}^{2}\left(s_{2}^{r}\right)^{4}-\sum_{r} \frac{E_{r} A_{r}}{b_{r}} z_{r}\left(s_{2}^{r}\right)^{2} s_{p}^{r} s_{q}^{r} \sum_{s} c_{p q}^{s} z_{s}\left(s_{2}^{s}\right)^{2}  \tag{4}\\
D_{12}= & v_{m} D+\sum_{r} \frac{E_{r} A_{r}}{b_{r}} z_{r}^{2}\left(s_{1}^{r}\right)^{2}\left(s_{2}^{r}\right)^{2}-\frac{1}{4} \sum_{r} \frac{E_{r} A_{r}}{b_{r}} z_{r}\left(s_{1}^{r}\right)^{2} s_{p}^{r} s_{q}^{r} \sum_{s} c_{p q}^{s} z_{s}\left(s_{2}^{s}\right)^{2}- \\
& -\frac{1}{4} \sum_{r} \frac{E_{r} A_{r}}{b_{r}} z_{r}\left(s_{2}^{r}\right)^{2} s_{p}^{r} s_{q}^{r} \sum_{s} c_{p q}^{s} z_{s}\left(s_{1}^{s}\right)^{2} \\
D_{66}= & \frac{\left(1-v_{m}\right) D}{2}+\sum_{r} \frac{E_{r} A_{r}}{b_{r}} z_{r}^{2}\left(s_{1}^{r}\right)^{2}\left(s_{2}^{r}\right)^{2}-\frac{1}{4} \sum_{r} \frac{E_{r} A_{r}}{b_{r}} z_{r} s_{1}^{r} s_{2}^{r} s_{p}^{r} s_{q}^{r} \sum_{s} c_{p q}^{s} z_{s} s_{1}^{s} s_{2}^{s}
\end{align*}
$$

$$
\begin{align*}
D_{26}= & \sum_{r} \frac{E_{r} A_{r}}{b_{r}} z_{r}^{2} s_{l}^{r}\left(s_{2}^{r}\right)^{3}-\frac{l}{4} \sum_{r} \frac{E_{r} A_{r}}{b_{r}} z_{r} s_{2}^{r} s_{2}^{r} s_{p}^{r} s_{q}^{r} \sum_{s} c_{p q}^{s} z_{s} s s_{1}^{s} s_{2}^{s}- \\
& -\frac{l}{4} \sum_{r} \frac{E_{r} A_{r}}{b_{r}} z_{r} s_{2}^{r} s_{1}^{r} s_{p}^{r} s_{q}^{r} \sum_{s} c_{p q}^{s} z_{s} s s_{2}^{s} s_{2}^{s}  \tag{4}\\
D_{16}= & \sum_{r} \frac{E_{r} A_{r}}{b_{r}} z_{r}^{2}\left(s_{l}^{r}\right)^{3} s_{2}^{r}-\frac{l}{4} \sum_{r} \frac{E_{r} A_{r}}{b_{r}} z_{r} s_{l}^{r} s_{l}^{r} s_{p}^{r} s_{q}^{r} \sum_{s} c_{p q}^{s} z_{s} s_{l}^{s} s_{2}^{s}- \\
& -\frac{l}{4} \sum_{r} \frac{E_{r} A_{r}}{b_{r}} z_{r} s_{1}^{r} s_{2}^{r} s_{p}^{r} s_{q}^{r} \sum_{s} c_{p q}^{s} z_{s} s s_{l}^{s} s_{l}^{s}
\end{align*}
$$

in which

$$
\begin{array}{ll}
D=\frac{E_{m} h^{3}}{12\left(1-v_{m}{ }^{2}\right)}, \quad \mathbf{c}^{r}=\operatorname{col}\left(c_{11}^{r}, c_{12}^{r}, c_{22}^{r}\right)=\frac{E_{r} A_{r}}{b_{r}} \mathbf{K}^{-1} \mathbf{s}^{r}, \quad \mathbf{K}=\left|k_{i j}\right|_{3 x 3} \\
k_{11}=\frac{E_{m} h}{1-v_{m}{ }^{2}}+\sum_{r} \frac{E_{r} A_{r}}{b_{r}}\left(s_{1}^{r}\right)^{4}, & k_{12}=2 \sum_{r} \frac{E_{r} A_{r}}{b_{r}}\left(s_{1}^{r}\right)^{3} s_{2}^{r} \\
k_{13}=\frac{E_{m} h v_{m}}{1-v_{m}{ }^{2}}+\sum_{r} \frac{E_{r} A_{r}}{b_{r}}\left(s_{1}^{r}\right)^{2}\left(s_{2}^{r}\right)^{2}, & k_{22}=2\left[\frac{E_{m} h}{1+v_{m}}+2 \sum_{r} \frac{E_{r} A_{r}}{b_{r}}\left(s_{1}^{r}\right)^{2}\left(s_{2}^{r}\right)^{2}\right] \\
k_{23}=2 \sum_{r} \frac{E_{r} A_{r}}{b_{r}}\left(s_{2}^{r}\right)^{3} s_{1}^{r}, & k_{33}=\frac{E_{m} h}{1-v_{m}{ }^{2}}+\sum_{r} \frac{E_{r} A_{r}}{b_{r}}\left(s_{2}^{r}\right)^{4}
\end{array}
$$

with notations: $E_{m}$ is the modulus of elasticity of the matrix, $\nu_{m}$-Poisson's ratio of the matrix, $h$ - thickness of the plate, $E_{r}$ - the modulus of elasticity of the $r$ th family of fibres, $A_{r}$ - cross-sectional area of an individual fibre, $z_{r}$ - the distance of the fibres of the $r$ th family from the middle plane of the plate, $b_{r}$ spacing of the $r$ th family of the fibres, $s_{l}^{r}$-cosine of angle between $x$-axis and axis of $r$ th family of fibres, $s_{2}^{r}$ - cosine of angle between $y$-axis and axis of $r$ th family of fibres.

### 2.2. The layerwise model

The stiffness coefficients $D_{i j}$ for the layerwise model are defined by

$$
\begin{equation*}
D_{i j}=\frac{1}{3} \sum_{k=1}^{N} \tilde{C}_{i j}^{(k)}\left(z_{k+1}^{3}-z_{k}^{3}\right) \tag{5}
\end{equation*}
$$

The cross-section of the analysed layered plate is shown in fig.2, where only two layers are marked for clarity. Location of the typical $k$ th layer and its thickness $h_{k}$ are defined by the coordinates $z_{k}$ and $z_{k+1}$. A typical fibre layer is shown in fig. 3. All such layers are assumed to be orthotropic in their material coordinate systems $x_{1} x_{2} x_{3}$. The global coordinate system, in which the corresponding boundary problem is solved, is denoted by $x y z$.


Fig. 2. Layered plate


Fig. 3. Typical $k$ th layer of fibre-composite plate, $k=1,2 \ldots N$
The constitutive equations between stress $\boldsymbol{\sigma}^{(k)}$ and strain $\boldsymbol{\varepsilon}^{(k)}$ tensors for the $k$ th orthotropic layer in the material coordinate system $x_{1} x_{2} x_{3}$ may be written as

$$
\begin{gathered}
\boldsymbol{\sigma}^{(k)}=\mathbf{C}^{(k)} \boldsymbol{\varepsilon}^{(k)} \\
{\left[\begin{array}{l}
\sigma_{11} \\
\sigma_{22} \\
\sigma_{12}
\end{array}\right]^{(k)}=\left[\begin{array}{ccc}
C_{11} & C_{12} & 0 \\
C_{12} & C_{22} & 0 \\
0 & 0 & C_{66}
\end{array}\right]^{(k)}\left[\begin{array}{c}
\varepsilon_{11} \\
\varepsilon_{22} \\
2 \varepsilon_{12}
\end{array}\right]^{(k)}}
\end{gathered}
$$

where

$$
\begin{aligned}
C_{11} & =\frac{E_{1}}{1-v_{12} v_{21}} \quad C_{22}=\frac{E_{2}}{1-v_{12} v_{21}} \\
C_{12} & =C_{21}=\frac{v_{12} E_{2}}{1-v_{12} v_{21}}=\frac{v_{21} E_{1}}{1-v_{12} v_{21}} \quad C_{66}=G_{12}
\end{aligned}
$$

$$
\begin{array}{rlr}
G_{r}=\frac{E_{r}}{2\left(l+v_{r}\right)} & G_{m}=\frac{E_{m}}{2\left(l+v_{m}\right)} \\
E_{l}=c_{r} E_{r}+c_{m} E_{m} & E_{2}=\frac{E_{r} E_{m}}{c_{r} E_{r}+c_{m} E_{m}} \\
v_{l 2}=c_{r} v_{r}+c_{m} v_{m} & G_{l 2}=\frac{G_{r} G_{m}}{c_{r} G_{r}+c_{m} G_{m}}
\end{array}
$$

Here: $E_{r}$ - the modulus of elasticity of $r$ th family of fibres (placed in $k$ th layer), $v_{r}$ - Poisson's ratio of $r$ th family of fibres, $c_{r}$ - volume fraction of $r$ th family of fibres, $E_{m}$ - the modulus of elasticity of matrix, $v_{m}$-Poisson's ratio of matrix, $c_{m}$ - volume fraction of matrix.

The orthotropic constitutive equations become anisotropic in the global coordinate system. The anisotropy is characterised by matrix $\widetilde{\mathbf{C}}^{(k)}$ and we have

$$
\left.\begin{array}{c}
\tilde{\boldsymbol{\sigma}}^{(k)}=\widetilde{\mathbf{C}}^{(k)} \tilde{\boldsymbol{\varepsilon}}^{(k)} \\
{\left[\begin{array}{l}
\boldsymbol{\sigma}_{x x} \\
\boldsymbol{\sigma}_{y y} \\
\boldsymbol{\sigma}_{x y}
\end{array}\right]^{(k)}=\left[\begin{array}{lll}
\tilde{C}_{11} & \tilde{C}_{12} & \tilde{C}_{16} \\
\tilde{C}_{12} & \tilde{C}_{22} & \tilde{C}_{26} \\
\tilde{C}_{16} & \tilde{C}_{26} & \tilde{C}_{66}
\end{array}\right]^{\boldsymbol{\varepsilon}_{x x}}{ }^{(k)}} \\
\boldsymbol{\varepsilon}_{y y} \\
2 \boldsymbol{\varepsilon}_{x y}
\end{array}\right]^{(k)}=\mathbf{T}^{(k)} \mathbf{C}^{(k)}\left(\mathbf{T}^{(k)}\right)^{T} .
$$

where $\mathbf{T}^{(k)}$ is the transformation matrix from the material coordinate system to the global coordinate system

$$
\mathbf{T}^{(k)}=\left[\begin{array}{ccc}
\cos ^{2} \theta & \sin ^{2} \theta & -2 \sin \theta \cos \theta \\
\sin ^{2} \theta & \cos ^{2} \theta & 2 \sin \theta \cos \theta \\
\sin \theta \cos \theta & -\sin \theta \cos \theta & \cos ^{2} \theta-\sin ^{2} \theta
\end{array}\right]
$$

with $\theta=\theta_{\mathrm{k}}$ being the angle between $x$ and $x_{l}$ for the $k$ th layer.

## 3. VARIATION FORMULATION

Equation (1) should be supplemented with appropriate boundary conditions. For the solution of (1) in practical problems one has to resort to approximate, numerical methods. We have used the finite element method which allows one to effectively solve complex boundary value problems (BVPs) of mechanics. To this aim we first formulate the BVP for (1) in a variational form. Let $U \subset H^{2}(\Omega)$ be a set of kinematically admissible displacements of the plate, $V \subset H^{2}(\Omega)$ a set of testing functions (virtual displacements). By $H^{2}(\Omega)$ we designate the Hilbert space of functions that are square integrable with up to their second derivative on $\Omega$. Further, let $\Gamma$ denote the boundary of $\Omega$ and $\Gamma=\Gamma_{D} \cup \Gamma_{N} \Gamma_{D} \cap \Gamma_{N}=0, \Gamma_{D}$ being the part of the boundary where constraints on displacements are imposed and on $\Gamma_{N}$-stress boundary conditions. The weak (variational) formulation of the BVP for (1) may be obtained from the principle of virtual work and stated as follows.

Find $w \in U$ such that

$$
\begin{equation*}
a(w, v)=f(v) \quad \text { for all } \quad v \in V \tag{6}
\end{equation*}
$$

The bilinear form $a(\bullet, \bullet)$ and the linear form $f(\bullet)$ are

$$
\begin{aligned}
& a(w, v)=\int_{\Omega}\left(D_{l 1} \frac{\partial^{2} w}{\partial x^{2}} \frac{\partial^{2} v}{\partial x^{2}}+2 D_{12} \frac{\partial^{2} w}{\partial x^{2}} \frac{\partial^{2} v}{\partial y^{2}}+4 D_{16} \frac{\partial^{2} w}{\partial x^{2}} \frac{\partial^{2} v}{\partial x \partial y}+\right. \\
&\left.+4 D_{26} \frac{\partial^{2} w}{\partial x \partial y} \frac{\partial^{2} v}{\partial y^{2}}+D_{22} \frac{\partial^{2} w}{\partial y^{2}} \frac{\partial^{2} v}{\partial y^{2}}+4 D_{66} \frac{\partial^{2} w}{\partial x \partial y} \frac{\partial^{2} v}{\partial x \partial y}\right) d x d y \\
& f(v)=\int_{\Omega} p v d x d y+\int_{I_{N}}\left(M_{n n} \frac{\partial v}{\partial n}+M_{n s} \frac{\partial v}{\partial s}+Q_{n} v\right) d s
\end{aligned}
$$

and correspond to the elastic strain energy and the virtual work of external loads, respectively.

## 4. FEM APPROXIMATION

In the finite element method the middsurface plane of the plate, $\Omega$, is divided into a number of finite elements $\Omega^{e} \subset \Omega$, and within each element $e$ the displacement field $w(x, y)$ is approximated by $w^{e}(x, y)$. Then, the displacement $w$ can be expressed as

$$
\begin{equation*}
w(x, y) \approx \sum_{e} w^{e}(x, y)=\sum_{e} \mathbf{N}^{e}(x, y) \mathbf{q}^{e}=\mathbf{N} \mathbf{q} \tag{7}
\end{equation*}
$$

where $\mathbf{N}^{e}=\mathbf{N}^{e}(x, y)$ is the matrix of shape functions and $\mathbf{q}^{\mathbf{e}}$ is the vector of the nodal parameters for element $e$, whereas $\mathbf{N}$ and $\mathbf{q}$ are their counterparts for the whole plate. Using (7) we can define stresses and strains in the plate as a function of $\mathbf{q}$ :

- field of strains

$$
\boldsymbol{\varepsilon}=\left[\begin{array}{l}
\varepsilon_{x x}  \tag{8}\\
\varepsilon_{y y} \\
\gamma_{x y}
\end{array}\right]=-z\left[\begin{array}{c}
w_{, x x} \\
w_{, y y} \\
2 w_{, x y}
\end{array}\right]=-z \boldsymbol{\kappa}=-z \mathbf{B} \mathbf{q}
$$

where $\boldsymbol{\kappa}$ stand for the matrix of curvatures

$$
\mathbf{\kappa}=\left[\begin{array}{c}
\kappa_{x x}  \tag{9}\\
\kappa_{y y} \\
2 \kappa_{x y}
\end{array}\right]=\mathbf{B} \mathbf{q}, \quad \mathbf{B}=\partial \mathbf{N}, \quad \partial=\left[\begin{array}{c}
\partial^{2} / \partial x^{2} \\
\partial^{2} / \partial y^{2} \\
2 \partial^{2} / \partial x \partial y
\end{array}\right]
$$

- field of stresses in $k t$ th layer

$$
\boldsymbol{\sigma}^{(k)}=\left[\begin{array}{l}
\sigma_{x x}  \tag{10}\\
\sigma_{y y}^{(k)} \\
\sigma_{x y}
\end{array}\right]^{=}=\tilde{\mathbf{C}}^{(k)} \boldsymbol{\varepsilon}^{(k)}=-z \tilde{\mathbf{C}}^{(k)} \mathbf{B} \mathbf{q}, \quad z \in\left[z_{k}, z_{k+l}\right]
$$

- field of moments

$$
\mathbf{M}=\left[\begin{array}{l}
M_{x x}  \tag{11}\\
M_{y y} \\
M_{x y}
\end{array}\right]=-\mathbf{D} \mathbf{B} \mathbf{q}
$$

The vector of nodal parameters $\mathbf{q}$ can be eventually determined from the system of the equilibrium equations

$$
\begin{equation*}
K q=f \tag{12}
\end{equation*}
$$

in which $\mathbf{K}$ is the global stiffness matrix and $\mathbf{f}$ is the global vector of nodal forces, both obtained by aggregation of all elemental contributions. In our own computer program we have used a 4-node quadrilateral finite element (PlQ4) with 12 degree of freedom (DOF). The explicit formulae for elements $K_{i j}$ of the stiffness matrix for this quadrilateral finite element are listed in [1]. For the sake of comparison, we have calculated test examples both with our program and the ABAQUS software.

## 5. NUMERICAL EXAMPLES

In this section we have presented and compared some of the results of numerical tests we have obtained for fibre-reinforced plates when using different models and various types of finite elements. The influence of finite element meshes is also considered.

## Exmaple 1.

At first we have analyzed a square simply supported orthotropic fibre-composite plate and compared the FEM results with the analytical Navier solution. The plate is reinforced by four families of fibres (see fig. 4). The thickness of the plate is 0.01 m and material coefficients are $E_{m}=3,15 \mathrm{GPa}, v_{m}=0,38, E_{r}=85,5$ GPa, $v_{r}=0,20$. The plate is loaded by a uniformly distributed pressure $q=1 \mathrm{kPa}$. We have solved the plate using the model of Świtka and the layerwise model, the corresponding stiffnesses $D_{i j}$ are gathered in table 1 . The used finite element was $20 \times 20$. The distribution of plate's deflection along half-span cross-section A-A for various cases is shown in fig.5. The solutions are similar with maximum differences of about $5 \%$.


Fig. 4. Orthotropic fibre-composite plate


Fig. 5. Deflection along the cross-section A-A in an orthotropic plate (cf. fig.4)

|  | $\mathrm{D}_{11}$ | $\mathrm{D}_{12}$ | $\mathrm{D}_{22}$ | $\mathrm{D}_{66}$ |
| :---: | :---: | :---: | :---: | :---: |
| Model of Świtka | $5.473 \mathrm{E}-4$ | $1.166 \mathrm{E}-4$ | $5.347 \mathrm{E}-4$ | $9.511 \mathrm{E}-5$ |
| Layerwise model | $5.541 \mathrm{E}-4$ | $1.209 \mathrm{E}-4$ | $5.426 \mathrm{E}-4$ | $1.07 \mathrm{E}-4$ |

Tab.1. Example 1: stiffness coefficients $D_{i j}$ by two models, with $\mathrm{D}_{16}=\mathrm{D}_{26}=0$

## Example 2.

Next, we have tested the fibre-composite plate that was composed of polyester matrix and eight families of glass $S$ fibres. Figure 6 shows cross-section of this plate. In our tests some selected families of fibres are rotated from the global coordinate system by angle $\theta_{i}$. The angle $\theta_{i}$ in the layerwise model is between global and material coordinate systems. Thickness of the plate is $h=0.02 \mathrm{~m}$ and material coefficients are $E_{m}=3,15 \mathrm{GPa}, v_{m}=0,38, E_{r}=85,5 \mathrm{GPa}, v_{r}=0,20$. We solved two cases, C 1 and C 2 , with different orientations of the fibre families (in all the cases the reinforcement is symmetric to the middsurface plane of the plate, M denotes a layer of matrix without fibres):

C1: $[\mathrm{M} / 0 / 90 / 0+\theta / 90+\theta / 0 / 90 / 0+\theta / 90+\theta / \mathrm{M}] \mathrm{s}$,
C2: $[\mathrm{M} / 0 / 90 / 22,5 / 112,5 / 45 / 135 / 67,5 / 157,5 / \mathrm{M}] \mathrm{s}$.
We have analyzed rectangular plate ( $1,00 \mathrm{~m} \times 1,50 \mathrm{~m}$ ) loaded by the uniformly distributed preassure $(\mathrm{p}=1 \mathrm{kPa})$ and concentrated force $(\mathrm{P}=1 \mathrm{kN})$ applied at the central point of the plate (point C). The plate is fixed on two edges and simply supported on one edge, see fig. 7. We have compared displacements in characteristic points $B$ and $C$ (fig.7) of the plate determined by own program and ABAQUS with different elements and meshes (figs. 8-11) using the layered model.


Fig. 6. Cross-section of the analysed plate and "cell" in layer with fibres


Fig.7. Examples 2 and 3: rectangular plate $1.00 \times 1.50$


Fig.8. Case $\mathrm{C} 1\left(\theta=0^{\circ}\right)$ : deflection at point C for different elements and meshes


Fig.9. Case $\mathrm{C} 1\left(\theta=0^{\circ}\right)$ : deflection at point B for different elements and meshes


Fig.10. Case C 2 : deflection at point C for different elements and meshes


Fig.11. Case C2: deflection at point B for different elements and meshes

## Example 3.

We have analyzed the same plate as in example 2, case 1 : $[\mathrm{M} / 0 / 90 / 0+\theta / 90+\theta /$ $0 / 90 / 0+\theta / 90+\theta / \mathrm{M}] \mathrm{s}$ and solved it for various values of angle $\theta$ in the indicated layers. Figures 12 to 14 show the influence of $\theta$ on the coefficients of stiffness $D_{i j}$ in both the models. Next, we analysed this influence on displacements (figs. 15-16) and stresses (figs. 17-20). It can be seen that although the differences in stiffnesses are about $13 \%$ in the case of $D_{11}$ and about $40 \%$ in $D_{12}$, this is not to such extent reflected in displacements whose changes are less then $5 \%$. Finally, figs. 17 and 18 illustrate distributions of shear stress $\sigma_{x y}$ for two values of $\theta$ by means of level sets.


Fig.12. Influence of fibre orientation on $D_{1 l}$ by different models


Fig.13. Influence of fibre orientation on $D_{12}$ by different models


Fig.14. Influence of fibre orientation on $D_{13}$ by different models


Fig.15. Case C 1 : influence of fibre orientation on deflection at point C by different models and programs


Fig.16. Case C 1 : influence of fibre orientation on deflection at point B by different models and programs


Fig.17. Case $\mathrm{C} 1\left(\theta=0^{\circ}\right)$ : Stress $\sigma_{x y}$ on the bottom surface of plate


Fig.18. Case $\mathrm{C} 1\left(\theta=45^{\circ}\right)$ : stress $\sigma_{x y}$ on the bottom surface of plate

## 6. CONCLUSION

In the paper we have presented the results of numerical simulation of the bending problem for a thin fibre-reinforced plate in the elastic range, making use of a number of models. The corresponding differential equation of the anisotropic plate is solved by the finite element method. In calculations we have used our own computer program and the program ABAQUS. The obtained results show some noticable differences in values of stiffness coefficients by the approach of Świtka and the layerwise approach, however, the resulting differences in displacements are of minor importance. If fact, the predictions of Świka's model are close to the results by ABAQUS with the rebar model. The applied 4-node quadrilateral plate element has proved to be quite effective, except for regions with point loads where the solution convergences in oppposite direction to that of ABAQUS. Finally, it may be concluded that the differences in the final quantities of interest due to the used different mathematical models of fiber-reinforcement are of minor importance from the enginnering point of view.

## Acknowledgement

Work supported by the Committee for Scientific Research (KBN) under Grant No 5 T07A 04224 in years 2003-2006. This support is gratefully acknowledged.

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# ANALIZA SPRĘŻYSTYCH PŁYT ZBROJONYCH WŁÓKNAMI 

## Streszczenie

W niniejszej pracy rozważa się zagadnienie zginania włóknokompozytowych płyt w zakresie sprężystym. Dla porównania, współczynniki sztywności płyty na zginanie i skręcanie wyznacza się według modelu Świtki i modelu warstwowego. Do wyznaczenia rozwiązania różniczkowego równania anizotropowej płyty zastosowano metodę elementtów skończonych. Rozwiązanie numeryczne wyznaczono za pomoca opracowanego własnego programu komputerowego i wykorzystujac program ABAQUS celem dokonania porównania. Praca zawiera wyniki wielu analiz numerycznych, wskazujące na wpływ kierunku ułożenia włókien oraz zastosowanych elementów skończonych i gęstości siatki elementów na wielkości fizyczne.

